

Will make effort to stay on syllabus numbering.

ECE 440: Lecture 12

Diffusion Current

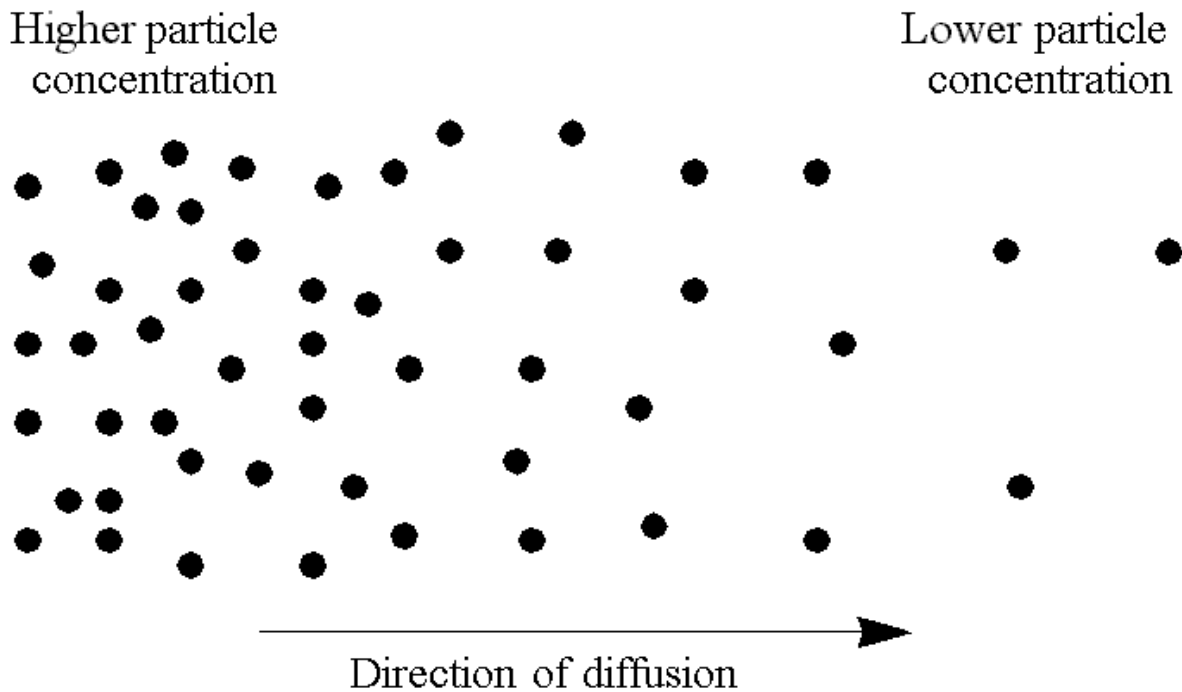
Remember Brownian motion of electrons & holes!

When E field = 0, but finite temperature ($T > 0$) thermal agitation velocity $v_T =$

But net drift velocity $v_d =$

So is net current $J_d =$ =

Now, what if there is a concentration gradient or a Brownian velocity gradient?

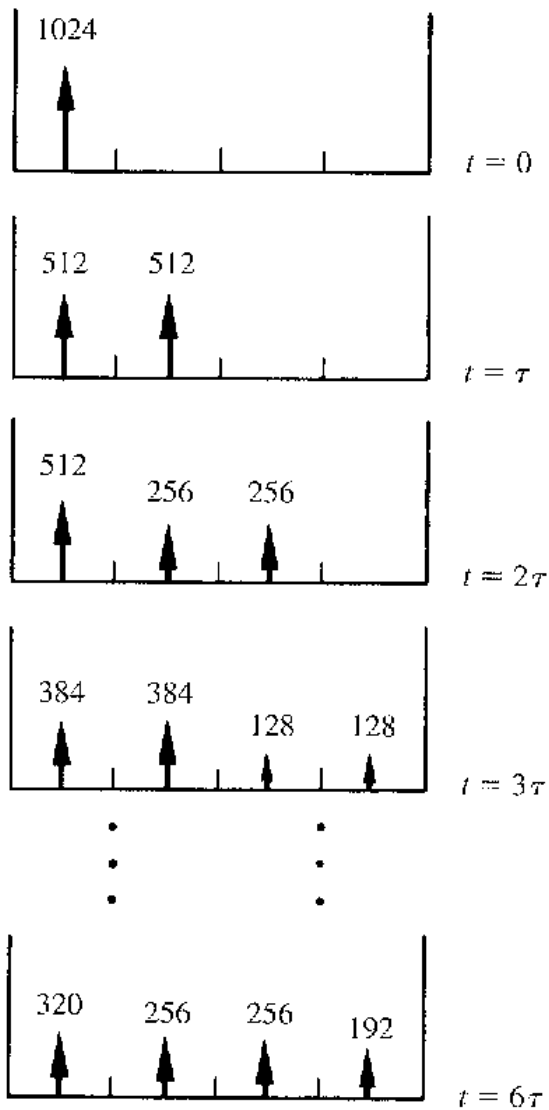


Is there a net flux of particles? Is there net current?

Examples of diffusion:

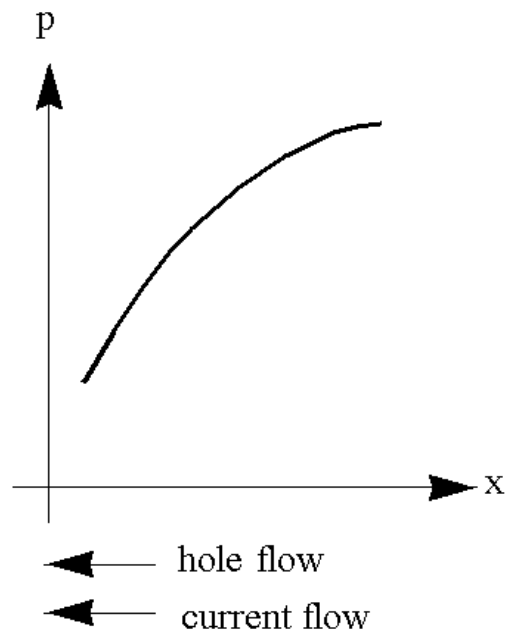
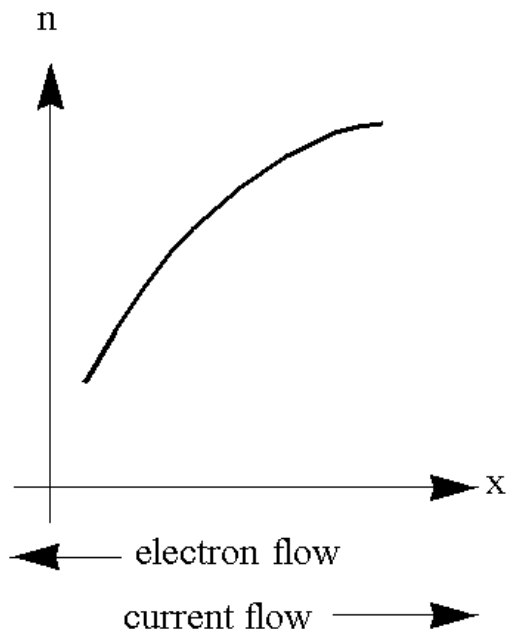
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One-dimensional diffusion example:



How would you set up diffusion in a semiconductor? You need something to drive it out of equilibrium.

What drives the *net* diffusion current? The concentration gradients! (no n or p gradients, no net current).



Mathematically:

$$\mathbf{J}_{N,diff} =$$

$$\mathbf{J}_{P,diff} =$$

Where D_N and D_P are the diffusion coefficients or diffusivity.

Now, we can FINALLY write down the TOTAL current in a semiconductor:

For electrons:

$$J_N = J_{N,drift} + J_{N,diff} = qn\mu_n\mathcal{E} + qD_N \frac{dn}{dx}$$

For holes:

$$J_P = J_{P,drift} + J_{P,diff} = qp\mu_p\mathcal{E} - qD_P \frac{dp}{dx}$$

And TOTAL current:

$$J =$$

Interesting point: minority carriers contribute little to drift current (usually, too few of them!), BUT if their gradient is high enough...

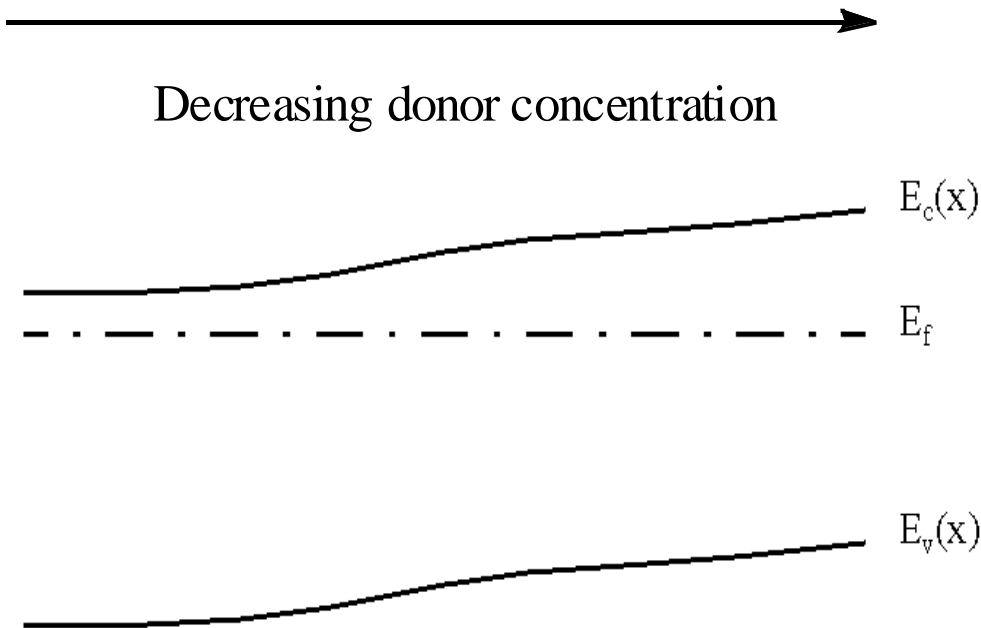
Under equilibrium, or open circuit conditions, total current must always be =

I.e. $J_{\text{drift}} = -J_{\text{diffusion}}$

More mathematically, for electrons:

So, any disturbance (e.g. light, doping gradient, thermal gradient) which may set up a carrier concentration gradient, will also internally set up a built-in _____

**What's the relationship between mobility and diffusivity?
Consider this band diagram (see, these things are *useful*):**



Going back to *Drift + Diffusion = 0* in equilibrium:

$$J_N = qn\mu_n \mathcal{E} + qD_N \frac{dn}{dx} = 0$$

Leads us to the Einstein Relationship:

$$\frac{D}{\mu} = \frac{kT}{q}$$

This is very, very important because it connects diffusivity with mobility, which we already know how to look up. Plus, it rhymes in many languages so it's easy to remember.

**The Einstein Relationship (almost) always holds true.
So we'll assume this in ECE 440.**

Example:

The hole density in an n-type silicon wafer ($N_d = 10^{17} \text{ cm}^{-3}$) decreases linearly from 10^{14} cm^{-3} to 10^{13} cm^{-3} between $x = 0$ and $x = 1 \text{ mm}$ (why?). Calculate the hole diffusion current.

PS: we've been surreptitiously neglecting something.

ECE 440: Lecture 13/14

Diffusion with Recombination

Last time:

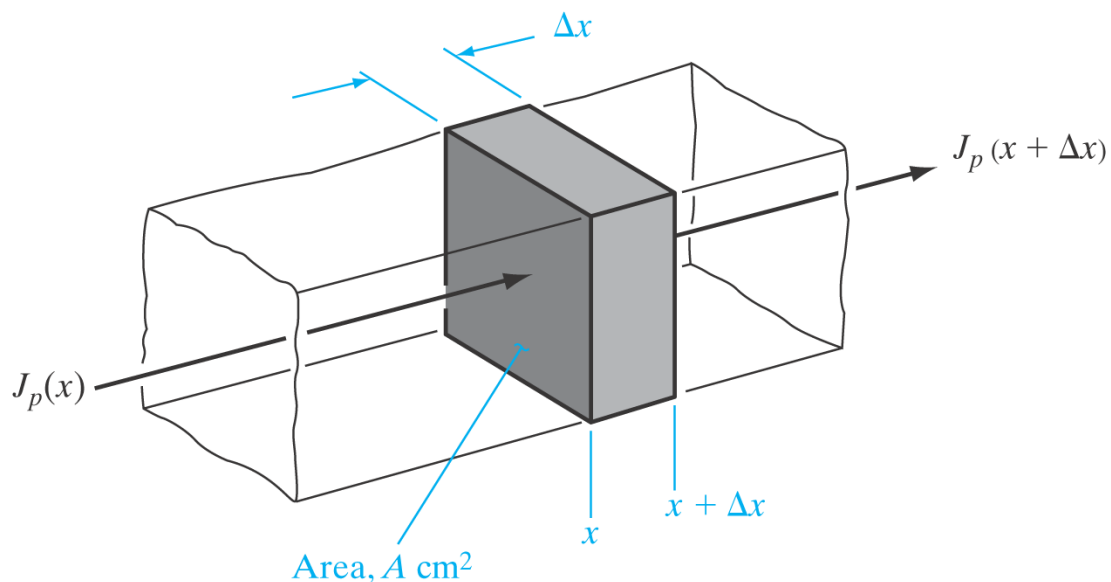
- **Diffusion without recombination (driven by dn/dx)**
- **Einstein relationship ($D/\mu = kT/q$)**
- **kT/q at room temperature ~ 0.026 V (be careful at temperatures different from 300 K)**
- **Mobility μ look up in tables, get diffusivity (be careful with total background doping concentration, N_A+N_D)**

Today:

- **Diffusion with recombination**
- **The diffusion length (distance until they recombine)**

Assume holes (p) are minority carriers.

Consider simple volume element where we have both generation, recombination, and holes passing through due to a concentration gradient (dp/dx).



This is simple “bean counting” in the little volume.

Rate of bean population increase = (current IN – current OUT) – bean recombination

Or replace “bean” with “bubble” in little fish tank volume.

Note, this technique is very powerful (and often used) in *any* Finite Element (FE) computational or mathematical model.

So, let's count beans (bubbles):

- **Recombination rate = #excess bubbles (δp) / recombination time (τ)**
- **Current (#bubbles) IN – Current (#bubbles) OUT = $J_{IN} - J_{OUT} / dx$**

Note units (VERY important check):

Bubble current is bubbles/cm²/second, but bubble rate of change (G&R) is bubbles/cm³/second.

So must account for width ($dx \sim cm$) of volume slice.

$$\frac{\partial p}{\partial t} = \frac{\partial \delta p}{\partial t} = -\frac{1}{q} \frac{\partial J}{\partial x} - \frac{\partial p}{\tau}$$

Why does the first equality hold? (it's simple, boring math)

Why is there a (diffusion) current derivative divided by q?

Of course, for (e.g.) holes,

$$J_{DIFF} = \quad \text{so,}$$

$$-\frac{1}{q} \frac{\partial J}{\partial x} =$$

So the diffusion equation (which is just a special case of the continuity equation above) becomes:

$$\frac{\partial \delta p}{\partial t} = D_P \frac{\partial^2 \delta p}{\partial x^2} - \frac{\delta p}{\tau}$$

This allows us to solve for the minority carrier concentration *in space and time* (here, holes).

Note, this is applicable only for minority carriers, whose net motion is dominated by diffusion.

What does this mean in *steady state*?

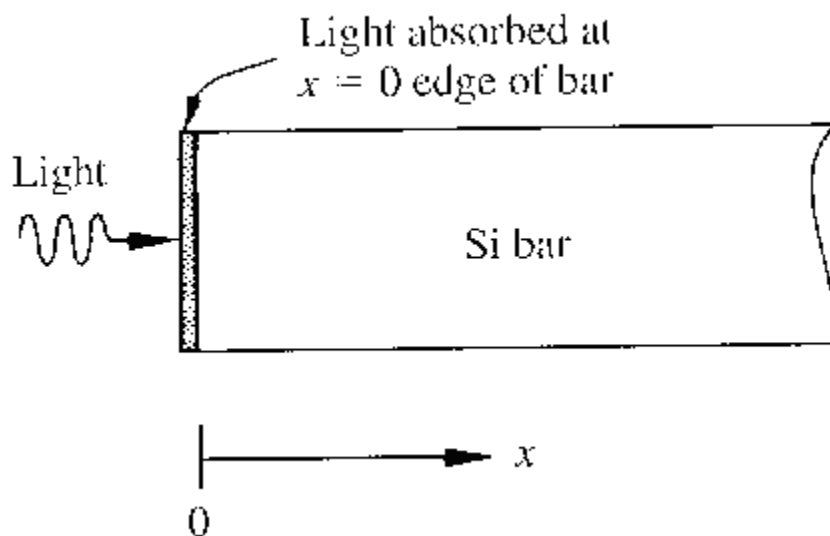
The diffusion equation in steady state:

$$\frac{\partial^2 \delta p}{\partial x^2} = \frac{\delta p}{D_p \tau} =$$

Interesting: this is what a lot of other diffusion problems look like in steady state. Other examples?

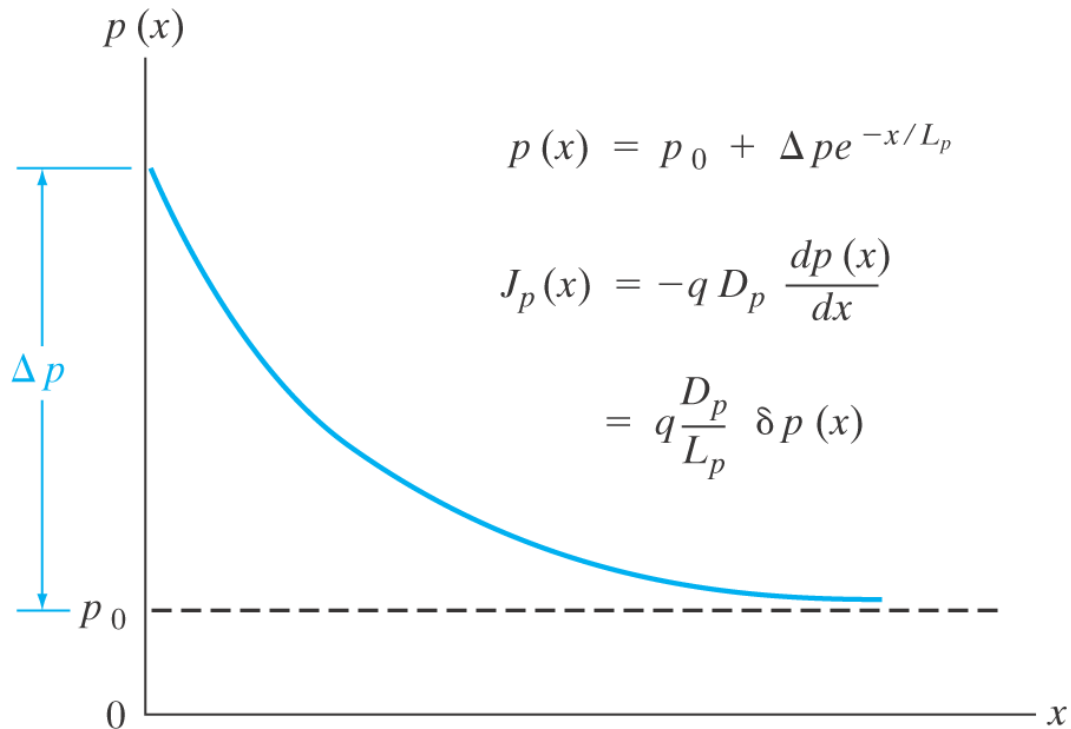
The *diffusion length* $L_p =$ is a figure of merit.

Consider an example under steady-state illumination:



Solve diffusion equation: $\delta p(x) = \Delta p e^{-x/L_p}$

Plot:



Physically, the diffusion lengths (L_P and L_N) are the average distance that minority carriers can diffuse into a sea of majority carriers before being annihilated (recombining).

What devices is this useful in?! (peek ahead)

Example: A) Calculate the minority carrier diffusion length in silicon with $N_D = 10^{16} \text{ cm}^{-3}$ and $\tau_p = 1 \text{ }\mu\text{s}$. B) Assuming 10^{15} cm^{-3} excess holes photogenerated at the surface, what is the diffusion current at $1 \text{ }\mu\text{m}$ depth?