

# Toward Low-Temperature Solid-Source Synthesis of Monolayer MoS<sub>2</sub>

Alvin Tang, Aravindh Kumar, Marc Jaikissoon, Krishna Saraswat, H.-S. Philip Wong, and Eric Pop\*

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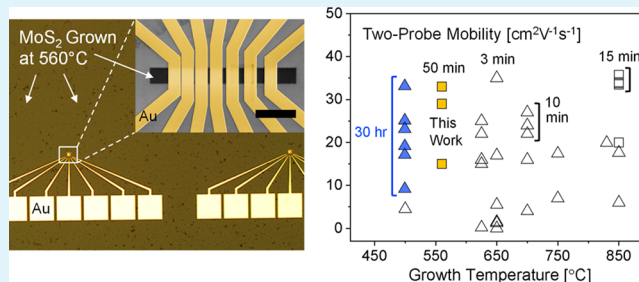
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**ABSTRACT:** Two-dimensional (2D) semiconductors have been proposed for heterogeneous integration with existing silicon technology; however, their chemical vapor deposition (CVD) growth temperatures are often too high. Here, we demonstrate direct CVD solid-source precursor synthesis of continuous monolayer (1L) MoS<sub>2</sub> films at 560 °C in 50 min, within the 450-to-600 °C, 2 h thermal budget window required for back-end-of-the-line compatibility with modern silicon technology. Transistor measurements reveal on-state current up to ~140 μA/μm at 1 V drain-to-source voltage for 100 nm channel lengths, the highest reported to date for 1L MoS<sub>2</sub> grown below 600 °C using solid-source precursors. The effective mobility from transfer length method test structures is 29 ± 5 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> at 6.1 × 10<sup>12</sup> cm<sup>-2</sup> electron density, which is comparable to mobilities reported from films grown at higher temperatures. The results of this work provide a path toward the realization of high-quality, thermal-budget-compatible 2D semiconductors for heterogeneous integration with silicon manufacturing.

**KEYWORDS:** 2D materials, transition metal dichalcogenides, MoS<sub>2</sub>, molybdenum disulfide, BEOL, back-end-of-the-line, chemical vapor deposition, CVD growth, carrier mobility



## INTRODUCTION

Since the era of the first integrated circuits, improving semiconductor device density has continuously translated into benefits for more advanced computing systems.<sup>1</sup> The horizontally stacked gate-all-around nanosheet transistor with 3 nm thick Si nanosheets is expected to replace the FinFET structure to continue transistor scaling beyond the 5 nm technology node.<sup>2</sup> However, further gate length scaling or better control over off-state leakage requires thinner body channel materials,<sup>3</sup> making atomically thin two-dimensional (2D) materials extremely appealing for use in next-generation computing devices. While graphene, the first 2D material discovered, is a semimetal with 0 eV band gap,<sup>4</sup> a class of 2D layered materials known as transition-metal dichalcogenides (TMDs) have band gaps ranging from semimetallic (0 eV) to semiconducting (~2.5 eV) energies.<sup>5,6</sup> TMDs retain sizable band gaps and carrier mobilities down to monolayer thicknesses<sup>7</sup> less than 1 nm, whereas Si thinned below ~4 nm faces severe mobility degradation issues due to surface roughness fluctuations.<sup>8–10</sup>

Specifically, MoS<sub>2</sub> has become one of the most promising TMDs due to its band gap (~2.2 eV for monolayer MoS<sub>2</sub><sup>11,12</sup>) and controllable growth of consistent large-area MoS<sub>2</sub> films down to monolayer (1L) thicknesses, making MoS<sub>2</sub> a viable candidate for large-scale semiconductor applications.<sup>13–18</sup>

Although moving MoS<sub>2</sub> away from tedious and inconsistent mechanical exfoliation techniques has enabled systematic studies for electronics applications,<sup>16,19,20</sup> the quality of synthesized MoS<sub>2</sub> has not always been ideal, especially from the standpoint of the device fabrication thermal budget. If MoS<sub>2</sub> is to be integrated in the back end (i.e., after silicon) of modern integrated circuits, the fabrication process of as-grown MoS<sub>2</sub> transistors cannot exceed the range of 450–600 °C (depending on the process time) for integration with logic applications,<sup>21–24</sup> although some three-dimensional (3D) vertical memory applications can withstand >700 °C.<sup>25</sup> For example, irreversible degradation of certain silicidic contacts, interlayer dielectrics, and diffusion barrier layers have been observed with thermal budgets over 600 °C for 2 h.<sup>21,26,27</sup>

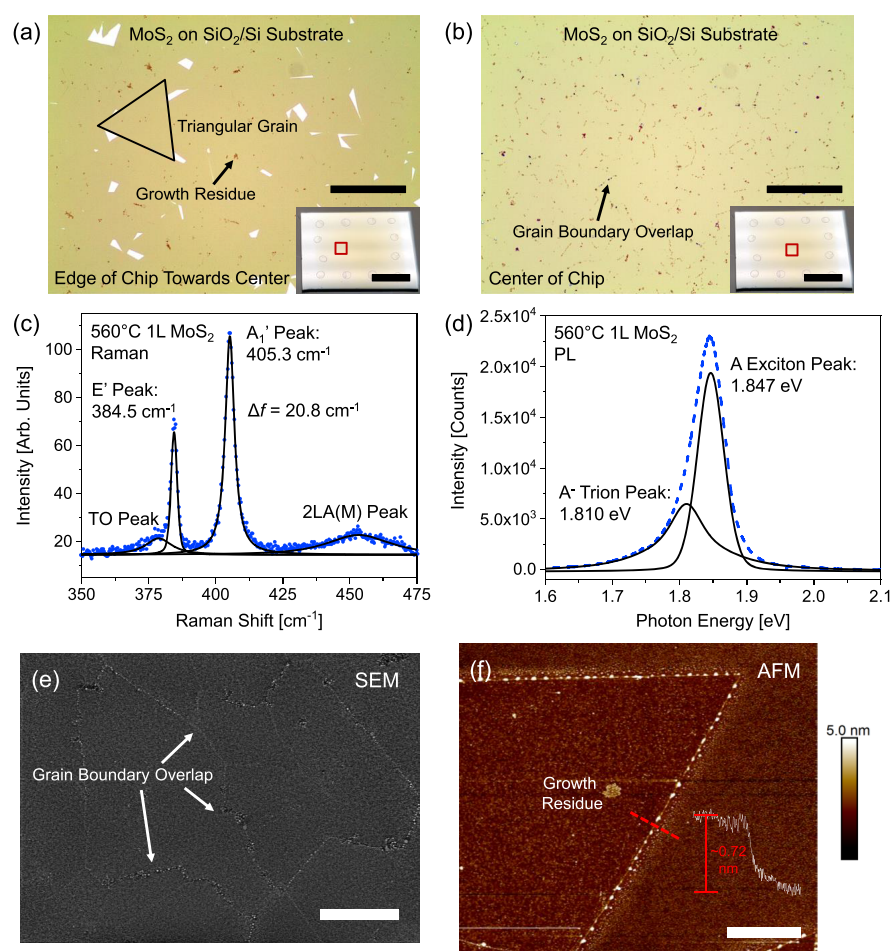
In this work, we report direct chemical vapor deposition (CVD) growth of continuous monolayer MoS<sub>2</sub> films using solid-source precursors at 560 °C in 50 min, within the thermal budget of modern integrated circuit processing. Comparing the

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**Figure 1.** Monolayer MoS<sub>2</sub> grown at 560 °C. (a,b) Optical images of monolayer (1L) MoS<sub>2</sub> grown by CVD at 560 °C on SiO<sub>2</sub>/Si at (a) 7 mm from the edge of the substrate and (b) the center of the substrate. Larger ~60 μm triangular grains are seen 7 mm from the edge, which merge into a continuous film with overlapping grain boundaries at the center. Scale bars are 20 μm. Insets show the 1.5 × 2 cm substrates, with red boxes marking the location of the optical images. Circles are “coffee rings” corresponding to the PTAS droplets. Inset scale bars are 7 mm. (c) Raman spectrum of 1L MoS<sub>2</sub> grown at 560 °C on SiO<sub>2</sub>/Si (blue dots). Pseudo-Voigt curves (black lines) fit the E' mode at 384.5 cm<sup>-1</sup> and the A<sub>1</sub>' mode at 405.3 cm<sup>-1</sup> with a peak separation Δf = 20.8 cm<sup>-1</sup>. The longitudinal acoustic 2LA(M) peak is centered at 453.3 cm<sup>-1</sup>, and the TO shoulder peak is centered at 378.8 cm<sup>-1</sup>. (d) PL spectrum of the same 1L MoS<sub>2</sub> (blue dashed line). Pseudo-Voigt curves (black lines) fit the A exciton peak at 1.847 eV and the charged A<sup>-</sup> trion peak at 1.810 eV. A laser with an excitation wavelength of 532 nm was used for all measurements. (e) SEM of the continuous 1L MoS<sub>2</sub> film at the center of the substrate with overlapping grain boundaries. Scale bar is 3 μm. (f) AFM of 1L MoS<sub>2</sub> triangular grains on SiO<sub>2</sub>/Si away from the substrate center. MoS<sub>2</sub> appears ~0.72 nm thick from the step height, in close agreement with the accepted monolayer MoS<sub>2</sub> thickness (0.615 nm).<sup>56</sup> Growth residue (most likely caused by excess partially reduced MoO<sub>3</sub> precursor) is observed on the MoS<sub>2</sub> grains and along the grain edges. Scale bar is 2 μm. Both SEM and AFM indicate that 1L MoS<sub>2</sub> is grown without bilayer (2L) MoS<sub>2</sub> nucleation.

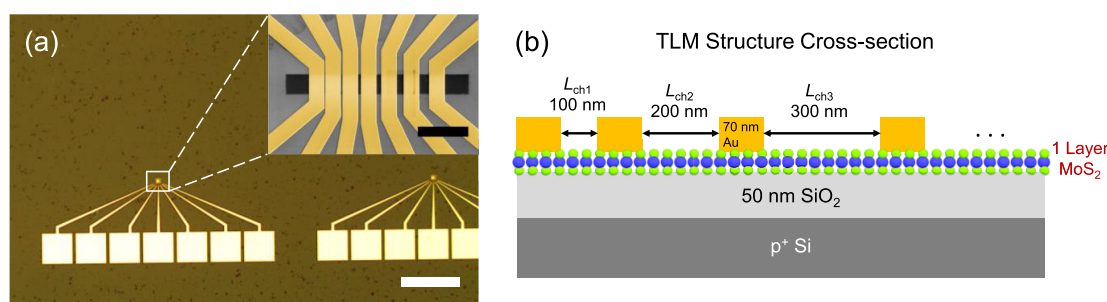
metrics of the  $I_{\text{on}}/I_{\text{off}}$  current ratio, drive current, contact resistance, and carrier mobility, transistors made from these as-grown films show performance levels similar to those reported for monolayer MoS<sub>2</sub> grown at higher temperatures.<sup>13,16,28–34</sup> This work provides a path toward the realization of high-quality, thermal-budget-compatible TMD materials, which could further improve future computing and data storage.

## ■ GROWTH AND MATERIAL CHARACTERIZATION

The thermal budget for modern silicon integrated circuit processing dictates that fabrication of BEOL (back-end-of-the-line) devices and materials cannot exceed 550 °C for 2 h,<sup>26,27</sup> and stricter temperature budgets are needed if process times are longer. This limitation is imposed by gate work function instability and silicide contact degradation.<sup>29</sup> Certain interlayer low- $\kappa$  dielectrics (e.g., SiCOH) could require even lower thermal budgets, below 450 °C for 2 h,<sup>21</sup> although SiCOH integrity could be maintained up to 525 °C for 2 h using some

BEOL processes.<sup>24</sup> Thus, in order to incorporate new materials such as MoS<sub>2</sub> into stacked, heterogeneous integrated circuits that increase the transistor density in the third dimension,<sup>35</sup> their BEOL thermal budget becomes crucial. However, direct monolayer (1L) MoS<sub>2</sub> growths with good electrical properties have typically been obtained using solid-source precursor CVD above 650 °C.<sup>13,16,28–34</sup> Wafer-scale 1L MoS<sub>2</sub> grown using metal–organic CVD (MOCVD)<sup>14,36–41</sup> at or below 500 °C appears within the BEOL temperature budget, but the growth times have been >8 h, which limit throughput, and the average carrier mobilities reported remain lower than the best MoS<sub>2</sub> grown at higher temperatures. Efforts to develop atomic layer deposition of MoS<sub>2</sub> at temperatures as low as 50 °C are also underway,<sup>42,43</sup> but the resulting films are amorphous and require post-growth anneals above 800 °C to improve their crystallinity.<sup>44–46</sup>

Here, we grew continuous 1L MoS<sub>2</sub> films at 560 °C in 50 min using solid sulfur (S) and molybdenum trioxide (MoO<sub>3</sub>)



**Figure 2.** TLM structures. (a) Optical image of the TLM structures with channels ranging from  $L_{\text{ch}} = 100$  to 700 nm. The large, square probe pads (20 nm  $\text{SiO}_2/2$  nm Ti/40 nm Au) lead into fine contacts (70 nm Au). Scale bar is 100  $\mu\text{m}$ . Inset: enlarged, colorized SEM of the fine leads directly contacting the monolayer  $\text{MoS}_2$  channel. The channel was etched 18  $\mu\text{m}$  long  $\times$  2  $\mu\text{m}$  wide (darker region) after the contact metal deposition. Inset scale bar is 5  $\mu\text{m}$ . (b) Cross-section schematic of the TLM test structure. The gold contact lengths are 1.5  $\mu\text{m}$  wide (not to scale). The Si substrate serves as the gate (G), and pairs of Au contacts serve as the drain (D) and grounded source (S) in subsequent measurements.

precursors with the aid of perylene-3,4,9,10 tetracarboxylic acid tetrapotassium salt (PTAS).<sup>13,15,17,18</sup> PTAS dissolved in water is deposited as droplets around the perimeter of 1.5  $\times$  2 cm chips with thermally grown  $\text{SiO}_2$  on  $\text{p}^+$  silicon, which also serves as the back-gate for field-effect transistors (FETs).<sup>13</sup> The PTAS droplets dry up into “coffee rings” around the perimeter of the substrate before the seeded  $\text{SiO}_2/\text{Si}$  chips are placed face down on top of an alumina ( $\text{Al}_2\text{O}_3$ ) crucible containing  $\text{MoO}_3$  powder. Trace amounts of PTAS diffuse from the dried “coffee rings” toward the center of the chip to facilitate  $\text{MoS}_2$  nucleation during growth. Excess solid sulfur source melted in a quartz boat is placed at an optimal position upstream near the incoming Ar gas flow inlet, all enclosed within a cylindrical quartz tube.

In addition to the PTAS seed layer promoting initial  $\text{MoS}_2$  nucleation from which grains grow, the trace amounts of carbon from the diffusing seed layer also act as a catalyst, promoting reduction of  $\text{MoO}_3$  powder into volatile suboxide  $\text{MoO}_{3-x}$ .<sup>47</sup> The reduction of  $\text{MoO}_3$  into  $\text{MoO}_{3-x}$  is crucial for incoming S to react with volatile molybdenum (Mo) species in order to grow  $\text{MoS}_2$ . It is important to note that at growth temperatures around 560  $^\circ\text{C}$ ,  $\text{MoO}_3$  reduction is near the limit of reaction with trace amounts of carbon present.<sup>47</sup> Both the reduced system pressure of 490 Torr and the incoming volatile sulfur introduced at elevated temperatures promote the reduction of  $\text{MoO}_3$  into volatile suboxide  $\text{MoO}_{3-x}$ .<sup>48</sup> When sufficient volatile  $\text{MoO}_{3-x}$  is matched with an optimal flux of incoming sulfur atoms on the substrate surface, further reaction occurs to grow 1L  $\text{MoS}_2$  films. As shown in Figure 1a,b, relatively clean 1L  $\text{MoS}_2$  films can be grown at 560  $^\circ\text{C}$  in 50 min on 1.5  $\times$  2 cm chips. This is the lowest thermal budget that produced clean, sizable (>10  $\mu\text{m}$ )  $\text{MoS}_2$  triangular grains merged into a continuous film using solid source precursors in our system. Larger  $\sim 60$   $\mu\text{m}$  triangular grains are seen (a) 7 mm from the edge of the substrate, which merge into a continuous film with overlapping grain boundaries at (b) the center of the substrate.

Raman spectroscopy and photoluminescence (PL) are utilized to confirm that 1L  $\text{MoS}_2$  is present from its characteristic phonon (lattice vibrations) and excitonic (electron–hole pair generation) properties, respectively. In order to ensure strong, well-defined detectable signals, all Raman and PL measurements were performed on 1L  $\text{MoS}_2$  grown on 90 nm  $\text{SiO}_2/\text{Si}$  substrates using a green laser with excitation wavelength of 532 nm. Due to the Lorentzian nature of the Raman (and PL) peaks and the Gaussian nature of the

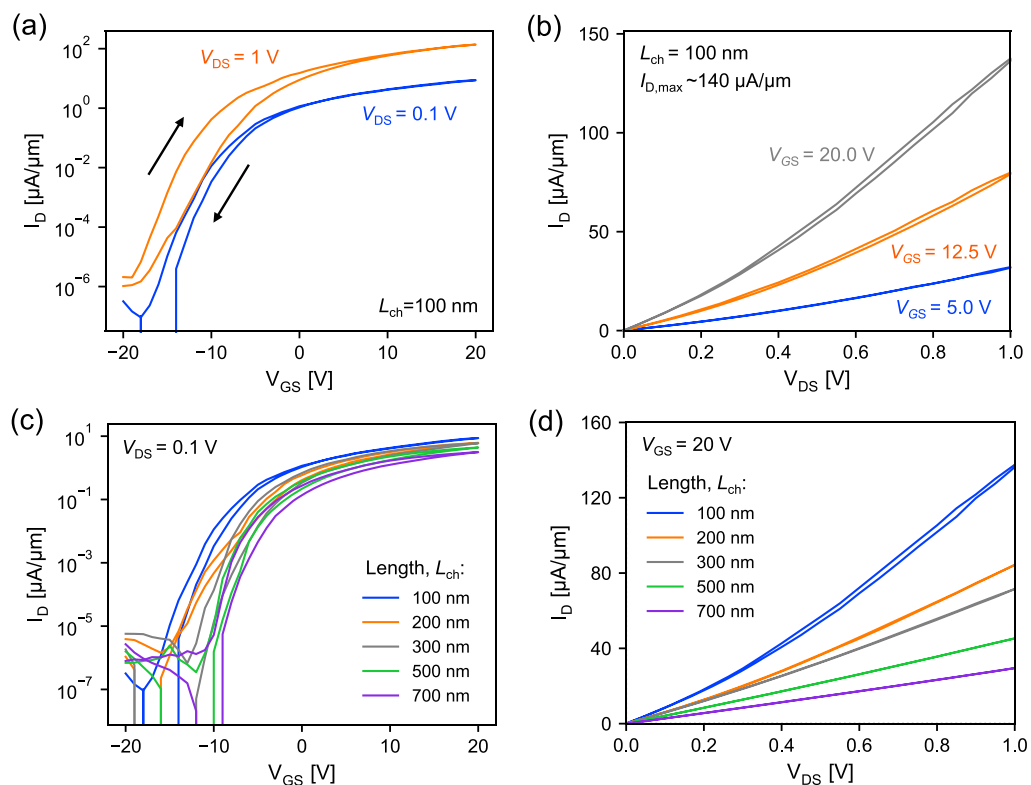
laser spatial intensity, pseudo-Voigt curves were used to capture each peak fit.<sup>49</sup> Figure 1c shows the Raman spectrum of 1L  $\text{MoS}_2$  grown at 560  $^\circ\text{C}$  (blue dots). Pseudo-Voigt curves (black lines) fit the Raman peaks of 1L  $\text{MoS}_2$ , namely, the in-plane  $E'$  mode at 384.5  $\text{cm}^{-1}$  and the out-of-plane  $A_1'$  mode at 405.3  $\text{cm}^{-1}$  with a peak separation  $\Delta f = 20.8$   $\text{cm}^{-1}$ , which is characteristic of as-grown 1L  $\text{MoS}_2$ <sup>13,50</sup> with slight intrinsic tensile strain due to thermal coefficient of expansion mismatch between the 2D material and the underlying  $\text{SiO}_2$  substrate.<sup>51</sup> Because odd-numbered, few-layer  $\text{MoS}_2$  (including 1L  $\text{MoS}_2$ ) belong to the  $D_{3h}$  point group, the two main Raman active modes are denoted  $E'$  and  $A_1'$ .<sup>13,52</sup> The 2LA(M) peak is seen around 453.3  $\text{cm}^{-1}$ , and the transverse optical (TO) shoulder peak is seen at 378.8  $\text{cm}^{-1}$ .<sup>53</sup>

Characteristic PL measurements are shown in Figure 1d. We note that these probe the optical band gap of 1L  $\text{MoS}_2$ , which is smaller than the electronic gap of 1L  $\text{MoS}_2$  by the exciton binding energy (0.2 to 0.6 eV),<sup>51</sup> and all three of these energies depend on the environmental dielectric screening.<sup>54</sup> Pseudo-Voigt curves (black lines) fit the PL spectrum, identifying the ground state A exciton at 1.847 eV and the charged  $A^-$  trion at 1.810 eV. (The  $A^-$  trion is a charged exciton with an extra electron coupled to the electron–hole pair.) We did not observe a B exciton in our samples (0.1 to 0.2 eV above the A peak), which has been used to assess nonradiative recombination, with low (or no) B peak suggesting higher sample quality.<sup>55</sup>

To further show that 1L  $\text{MoS}_2$  is grown by CVD at 560  $^\circ\text{C}$ , scanning electron microscopy (SEM) is shown in Figure 1e of the continuous 1L  $\text{MoS}_2$  film at the center of the substrate, with overlapping grain boundaries. Atomic force microscopy (AFM) of 1L  $\text{MoS}_2$  triangular grains on  $\text{SiO}_2/\text{Si}$ , away from the substrate center, is also shown in Figure 1f.  $\text{MoS}_2$  is measured to be  $\sim 0.72$  nm thick from the step height, which is in good agreement with the accepted monolayer  $\text{MoS}_2$  thickness (0.615 nm).<sup>56</sup> Growth residue (most likely caused by excess partially reduced  $\text{MoO}_3$  precursor) is observed on the  $\text{MoS}_2$  grains and along the grain edges. Both SEM and AFM indicate that 1L  $\text{MoS}_2$  is grown without bilayer (2L)  $\text{MoS}_2$  nucleation.

## ■ ELECTRICAL RESULTS AND DISCUSSION

Transfer length method (TLM) structures<sup>57</sup> were fabricated on 560  $^\circ\text{C}$  1L  $\text{MoS}_2$  as-grown on a 50 nm  $\text{SiO}_2$  on  $\text{p}^+$  Si substrate, which also serves as the back-gate for all devices. Figure 2a shows an optical image of the TLM structures as



**Figure 3.** Electrical characteristics. Transistors fabricated using monolayer MoS<sub>2</sub> grown at 560 °C on SiO<sub>2</sub>/Si substrates. (a) Measured  $I_D$  vs  $V_{GS}$  transfer curves at  $V_{DS} = 0.1$  and 1 V for a 100 nm channel device ( $L_{ch} = 100$  nm). The small arrows mark forward (left  $I_D$  branch) and backward (right  $I_D$  branch) voltage sweep directions, illustrating small hysteresis. A current ratio  $I_{on}/I_{off} \sim 10^7$  is observed, and the SS is estimated to be  $SS \approx 1150$  mV/dec from the forward sweep at  $V_{DS} = 0.1$  V. (b) Measured  $I_D$  vs  $V_{DS}$  curves at  $V_{GS} = 5$  to 20 V for the same  $L_{ch} = 100$  nm channel length device. The maximum drive current achieved was  $I_{D,max} \sim 140 \mu A/\mu m$  at  $V_{DS} = 1$  V. Current saturation was not observed because the threshold voltage  $V_T$  keeps this device in the linear operating region throughout.<sup>30</sup> (c) Measured  $I_D$  vs  $V_{GS}$  transfer curves at  $V_{DS} = 0.1$  V for transistors with channel lengths from  $L_{ch} = 100$  to 700 nm. (d) Measured  $I_D$  vs  $V_{DS}$  at  $V_{GS} = 20$  V for transistors with channel lengths from  $L_{ch} = 100$  to 700 nm.  $I_D$  decreases for increasing channel lengths as expected due to the larger channel resistance. All device channels are 2  $\mu m$  wide, and all measurements were performed at room temperature in a vacuum probe station. All electrical measurements shown are double sweeps (voltage swept up and down), revealing minimal hysteresis.

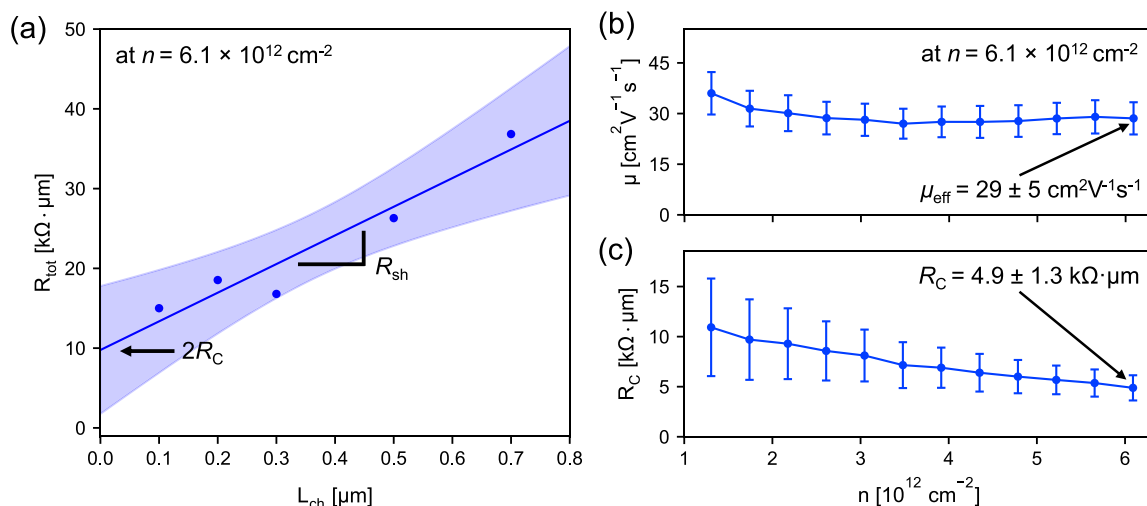
fabricated by electron-beam lithography. The large, square probe pads (20 nm SiO<sub>2</sub>/2 nm Ti/40 nm Au) lead into the fine contacts (70 nm Au). The colorized SEM inset shows the pure Au leads directly contacting the monolayer MoS<sub>2</sub> channel, which was etched 18  $\mu m$  long  $\times$  2  $\mu m$  wide (darker region) after the contact metal deposition. Figure 2b shows the cross-section schematic of the TLM test structure with adjacent contact channel lengths ranging from  $L_{ch} = 100$  to 700 nm, forming a series of back-gated FETs.

Figure 3a shows the measured drain current  $I_D$  (normalized by width) versus back-gate voltage  $V_{GS}$  transfer curves at drain voltages  $V_{DS} = 0.1$  and 1 V for a 100 nm channel device ( $L_{ch} = 100$  nm). The small arrows mark forward (left  $I_D$  branch) and backward (right  $I_D$  branch) voltage sweep directions, illustrating small hysteresis and charge trapping between MoS<sub>2</sub> and the back-gate SiO<sub>2</sub> dielectric. The maximum gate leakage observed was  $\sim 1$  nA at  $V_{GS} = 20$  V, which is 3–4 orders of magnitude smaller than the drain current in the on-state. A current ratio  $I_{on}/I_{off} \sim 10^7$  is observed, and the subthreshold slope (SS) is estimated  $SS \approx 1150$  mV/dec from the forward sweep at  $V_{DS} = 0.1$  V. We note the SS is relatively high but can be reduced by decreasing the equivalent oxide thickness of the gate dielectric below the 50 nm SiO<sub>2</sub> used here. Reducing interface charge trap density at the MoS<sub>2</sub>/dielectric interface<sup>58</sup> by passivating the dielectric<sup>59</sup> or by using

a molecular crystal seeding layer<sup>60</sup> can also greatly improve the SS when employing cleaner, dedicated fabrication outside general-purpose academic facilities.

Figure 3b shows the measured  $I_D$  versus  $V_{DS}$  output curves at back-gate voltages from  $V_{GS} = 5$  to 20 V for the same  $L_{ch} = 100$  nm device. The maximum drive current achieved was  $I_{D,max} \sim 140 \mu A/\mu m$  at  $V_{DS} = 1$  V and  $V_{GS} = 20$  V, which is the highest reported for monolayer MoS<sub>2</sub> grown below 600 °C. Current saturation is not observed because the threshold voltage  $V_T$  is sufficiently small to keep this device in the linear operating region throughout. Relatively low hysteresis is again observed from the forward and backward voltage sweeps, indicating minimal charge trapping. The measured  $I_D$  versus  $V_{GS}$  transfer curves at  $V_{DS} = 0.1$  V for channel lengths from  $L_{ch} = 100$  to 700 nm are shown in Figure 3c. The threshold voltage  $V_T$  ranges from 2.4 to 5.8 V using the linear extrapolation method.<sup>16</sup> The measured  $I_D$  versus  $V_{DS}$  output curves at  $V_{GS} = 20$  V for channel lengths from  $L_{ch} = 100$  to 700 nm are shown in Figure 3d.  $I_D$  decreases for increasing channel lengths, as expected, due to larger channel resistance contribution for a constant  $V_{DS}$  and a constant contact resistance ( $R_C$ ).

We next turn to estimating  $R_C$  first and then the electron mobility from TLM structures. In addition to a better estimation of  $R_C$ , the MoS<sub>2</sub> channel sheet resistance ( $R_{sh}$ ) can also be accurately extracted from the slope of the linear

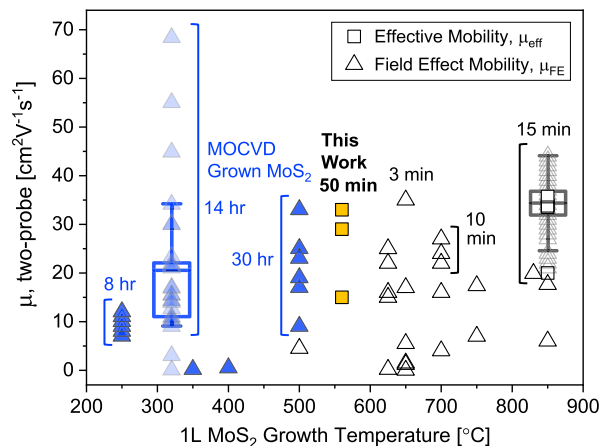


**Figure 4.** Sheet resistance and mobility. (a) Linear TLM fit of the total measured resistance ( $R_{\text{tot}}$ , normalized by channel width) as a function of the channel length ( $L_{\text{ch}}$ ) at a carrier concentration  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$ , corresponding to the same gate overdrive,  $V_{\text{OV}} = V_{\text{GS}} - V_{\text{T}} = 14 \text{ V}$ . The y-intercept corresponds to twice the contact resistance ( $2R_{\text{C}}$ ), while the slope of the line is the sheet resistance ( $R_{\text{sh}}$ ). The shaded region marks the 95% confidence bound of the line fit. (b) Effective mobility ( $\mu_{\text{eff}}$ ) vs  $n$  as obtained from the sheet resistance. Here,  $\mu_{\text{eff}} = 29 \pm 5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  at  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$ , which is comparable to previous reports of monolayer MoS<sub>2</sub> grown at higher temperatures (see Figure 5). (c) Contact resistance ( $R_{\text{C}}$ , normalized by width) as a function of the carrier concentration,  $n$ . The contact resistance was down to  $R_{\text{C}} = 4.9 \pm 1.3 \text{ k}\Omega \cdot \mu\text{m}$  at  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$ . All measurements were taken at  $V_{\text{DS}} = 0.1 \text{ V}$  and room temperature in a vacuum probe station.

TLM extrapolation. The sheet resistance is then used to estimate the effective mobility,  $\mu_{\text{eff}} = (qnR_{\text{sh}})^{-1}$ , where  $q$  is the elementary charge and  $n$  is the electron density per unit area. Here,  $n = C_{\text{ox}}(V_{\text{GS}} - V_{\text{T}} - V_{\text{DS}}/2)/q$  as all our transistors remain in the linear regime, where  $C_{\text{ox}}$  is the gate dielectric capacitance per unit area ( $\sim 70 \text{ nF/cm}^2$  for our 50 nm SiO<sub>2</sub>). We note  $V_{\text{T}}$  must be individually assessed for each channel in the TLM to account for any device-to-device variation.

Figure 4a shows the linear TLM fit of the total measured resistance ( $R_{\text{tot}} = L_{\text{ch}}R_{\text{sh}} + 2R_{\text{C}}$ , normalized by width) versus channel length ( $L_{\text{ch}}$ ) extracted from Figure 3d at a carrier concentration  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$  (at the same gate overdrive,  $V_{\text{OV}} = V_{\text{GS}} - V_{\text{T}} = 14 \text{ V}$ , not the same gate voltage  $V_{\text{GS}}$ ).  $V_{\text{T}}$  is extracted for each channel length in Figure 3c by the linear extrapolation method.<sup>57,61</sup> The y-intercept of the linear fit in Figure 4a corresponds to the total contact resistance,  $2R_{\text{C}}$ . We note that it is important to perform such TLM fits using a wide range of channel lengths (from short, dominated by  $2R_{\text{C}}$ , to long, dominated by  $L_{\text{ch}}R_{\text{sh}}$ ) in order to minimize the  $R_{\text{C}}$  extrapolation error. The remaining uncertainty of the linear TLM fit represents the intrinsic device-to-device variation. Figure 4b shows the effective mobility ( $\mu_{\text{eff}}$ ) as a function of carrier concentration  $n$ , obtained from the TLM plot slope, which is the sheet resistance ( $R_{\text{sh}}$ ). At  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$ , the effective electron mobility is  $\mu_{\text{eff}} = 29 \pm 5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , which is comparable to mobilities reported at higher growth temperatures.<sup>13,16,28–31,33,34</sup> Figure 4c shows the contact resistance ( $R_{\text{C}}$ , normalized by width) also as a function of the carrier concentration  $n$ . The contact resistance was  $R_{\text{C}} = 4.9 \pm 1.3 \text{ k}\Omega \cdot \mu\text{m}$  at  $n = 6.1 \times 10^{12} \text{ cm}^{-2}$  with the error bound corresponding to the 95% confidence interval of the line fits.

In Figure 5 we compare the room-temperature electron mobility data of our films grown at 560 °C with previous reports for other CVD-grown (and MOCVD-grown<sup>36–41</sup>) 1L MoS<sub>2</sub> reported in the literature,<sup>13,16,18,20,28–30,32,34,62–67</sup> as a function of the growth temperature. We note that some reports are given as effective mobility<sup>13,16,28–30</sup> ( $\mu_{\text{eff}}$  from TLM, denoted by squares); others are only available as field-effect



**Figure 5.** Monolayer MoS<sub>2</sub> electron mobility vs growth temperature. Reported effective mobilities ( $\mu_{\text{eff}}$  from TLM) are squares and two-probe field-effect mobilities ( $\mu_{\text{FE}}$ ) are triangles.  $\mu_{\text{eff}}$  from this report are marked in yellow. MOCVD data<sup>36–41</sup> are  $\mu_{\text{FE}}$  marked in blue, and solid-source CVD data from the literature<sup>13,16,18,20,28–30,32,34,62–67</sup> are hollow symbols. Growth times (at maximum process temperature) are also labeled, indicating significantly longer growths reported for MOCVD to attain full coverage 1L MoS<sub>2</sub>. The CVD films in this work were grown at 560 °C for 50 min, while MOCVD films were grown at or below 500 °C for 8 to 30 h,<sup>36–41</sup> as labeled. The  $\mu_{\text{eff}}$  from sheet resistance (TLM) measurements tend to be more reliable, while  $\mu_{\text{FE}}$  could be under- or overestimated depending on the  $V_{\text{GS}}$ -dependence of gated contacts.<sup>68,73,74</sup> All mobility data are reported at room temperature. During review, we became aware of recent MOCVD films grown at 320 °C for over 14 h;<sup>37,38</sup> a summary of  $\mu_{\text{FE}}$  data for these films (lighter blue triangles) is shown as the box-and-whisker plot at 320 °C (average  $\mu_{\text{FE}}$  of 20.4  $\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  in the blue box).

mobility<sup>18,20,32,34,36,62–67</sup> ( $\mu_{\text{FE}}$ , denoted by triangles). The effective mobility values from this work (grown at 560 °C for 50 min) are marked in yellow. The field-effect mobility values reported from MOCVD growths (at or below 500 °C for 8 to 30 h) are marked in blue.<sup>36–41</sup> For this simple comparison, we

benchmark two-probe mobility measurements, although more complex four-probe measurements with a threshold voltage correction can yield more accurate mobility values<sup>68</sup> if current shunt paths through the (invasive) inner voltage probes are avoided.<sup>7</sup> In general, we have found no correlation between carrier mobility and growth temperature, only between growth temperature and MoS<sub>2</sub> crystallite size or substrate adhesion. Growth times (at maximum process temperature) are also labeled, indicating significantly longer growths reported for MOCVD to attain full coverage 1L MoS<sub>2</sub>.

First, comparing the effective 1L MoS<sub>2</sub> mobilities ( $\mu_{\text{eff}}$  squares), we observe that our values (29 and 33 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> from 560 °C growths at a reduced 490 Torr pressure) are similar to the  $\mu_{\text{eff}}$  from MoS<sub>2</sub> grown at higher temperatures, up to 850 °C.<sup>13,16,28–30</sup> (Recent simulations have shown that mobilities in this range are limited by point defects, most likely, charge impurities and sulfur vacancies.<sup>69,70</sup>) Field-effect mobility ( $\mu_{\text{FE}}$ ) data reported for 1L MoS<sub>2</sub> grown at 850 °C (lighter hollow triangles) are also included as a box-and-whisker plot (average  $\mu_{\text{FE}}$  of 34.2 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> outlined in the black box),<sup>16</sup> in good agreement with the 850 °C effective mobility. For CVD-grown 1L MoS<sub>2</sub> using solid-source precursors, our values are the highest reported mobilities to date with a thermal budget below 2 h at 600 °C. One notable difference is that our past 850 °C growths have yielded up to ~10% bilayer regions,<sup>16</sup> while such bilayer regions are undiscernible on the 560 °C growths presented here (Figure 1). In addition, we have also found the MoS<sub>2</sub> grown at 560 °C to be more weakly adhered to the SiO<sub>2</sub> substrate compared to our previous studies at higher growth temperatures,<sup>13,16,28–30</sup> which necessitated careful electron-beam lithography without exposing the MoS<sub>2</sub> samples to water to avoid delamination during processing.<sup>71,72</sup> We note that adhesion challenges, in addition to the growth residue observed on the 1L MoS<sub>2</sub> grains and along the grain edges, could have also led to some of our observed device-to-device variation.

When comparing effective mobility values with field-effect mobilities ( $\mu_{\text{FE}}$ , triangles) reported across the literature, most  $\mu_{\text{FE}}$  reported are lower. This is mainly due to the effects of the contact resistance, although material quality, nonideal contact selection, and device fabrication (in academic facilities) could also cause the large  $\mu_{\text{FE}}$  variation observed, with no clear trends between studies and across growth temperatures. It is also possible that  $\mu_{\text{FE}}$  is overestimated in some studies depending on the gate-voltage-dependence of the contacts. In other words, a sharp turn-on of back-gated contacts can lead to an overestimated peak transconductance  $g_m$  and  $\mu_{\text{FE}}$ .<sup>68,73,74</sup> Due to these effects of gated contacts, effective mobility ( $\mu_{\text{eff}}$ ) from sheet resistance (TLM) measurements (particularly when reported as a function of the carrier density  $n$ ) tend to be more reliable and are preferable, rather than  $\mu_{\text{FE}}$ , as the main figure of merit for MoS<sub>2</sub>.

## CONCLUSIONS

We have demonstrated direct CVD solid source precursor growth of monolayer (1L) MoS<sub>2</sub> at 560 °C in 50 min, with electrical and optical properties very similar to those of CVD MoS<sub>2</sub> grown at higher temperatures. Our films are within the 450-to-600 °C, 2 h thermal budget window required for BEOL compatibility with modern silicon integrated circuit processing. These new 1L MoS<sub>2</sub> films growths were enabled by carefully matching an optimized sulfur flux to sufficient volatile MoO<sub>3-x</sub> reduced in the presence of the carbon-based PTAS seed layer.

Electrical measurements revealed a maximum drive current  $I_{\text{D,max}} \sim 140 \mu\text{A}/\mu\text{m}$  at  $V_{\text{DS}} = 1 \text{ V}$ , which is the highest reported for 1L MoS<sub>2</sub> grown below 600 °C using solid-source precursors. The effective electron mobility was extracted with TLM test structures as  $\mu_{\text{eff}} = 29 \pm 5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  at a carrier concentration of  $6.1 \times 10^{12} \text{ cm}^{-2}$ , which is comparable to mobilities reported from films grown at higher temperatures. The results of this work provide a path toward the realization of high-quality, thermal-budget-compatible TMDs for heterogeneous integration with silicon manufacturing. These could enable 3D integration of such 2D materials for advanced functionalities in memory or power-gating circuits (with low-leakage BEOL transistors) or high-density logic in the third dimension.

## METHODS

**560 °C 1L MoS<sub>2</sub> CVD Growth.** PTAS dissolved in water was deposited as droplets around the perimeter of  $1.5 \times 2 \text{ cm}$  chips of thermally grown 50 nm SiO<sub>2</sub> on p<sup>+</sup> silicon that were initially rinsed with deionized water.<sup>13</sup> The PTAS droplets dried into “coffee rings” around the perimeter of the substrate before the seeded SiO<sub>2</sub>/Si chips were placed face-down on top of an alumina (Al<sub>2</sub>O<sub>3</sub>) crucible containing ~2 mg of MoO<sub>3</sub> powder (Alfa Aesar, Puratronic 99.9995% purity) spread into a ~1 cm diameter circle. ~100 mg of excess solid sulfur source (Alfa Aesar, Puratronic 99.999% purity) melted in a quartz boat was placed ~26 cm away from the center near the incoming Ar gas flow inlet, all enclosed within a 55 mm inner diameter quartz tube. Before each growth, the system was flushed for 5 min using 1500 sccm Ar gas under vacuum before setting the ambient condition to 490 Torr at 22 sccm Ar flow rate, adjusting the throttle valve. For the growth temperature cycle, the tube furnace was first ramped from room temperature to 450 °C in 10 min and then subsequently ramped to 560 °C in 5 min. The temperature was held at 560 °C for 50 min before rapidly cooling down to room temperature by opening the furnace hatch. The partially reduced MoO<sub>3</sub> source must be cleaned out and replenished after each growth cycle, while the excess solid sulfur source can be reused in subsequent growths.

**MoS<sub>2</sub> Device Fabrication.** Contact probe pads and coarse contacts were defined by e-beam lithography using a bilayer resist stack of 250 nm methyl methacrylate (MMA)/50 nm poly (methyl methacrylate) (PMMA). This promotes resist undercutting for easier contact metal liftoff which helps maintain MoS<sub>2</sub> adhesion to the growth substrate. This was followed by e-beam evaporation of 20 nm SiO<sub>2</sub>/2 nm Ti/70 nm Au, all layers deposited at 0.5 Å/s at a base pressure of 10<sup>-7</sup> Torr. The 20 nm SiO<sub>2</sub> layer helps mitigate pad-to-gate leakage through the back-gate dielectric. This was followed by an overnight liftoff in acetone at room temperature. Fine contacts were patterned using e-beam lithography with a 60 nm PMMA 495k/215 nm PMMA 950k bilayer resist stack which provides fine resolution along with sufficient undercutting for effective liftoff. 70 nm Au was then deposited by e-beam evaporation at a deposition rate of 0.5 Å/s at a base pressure below 10<sup>-7</sup> Torr without any adhesion layer to achieve a clean contact interface. The liftoff was carried out in Remover PG solvent for 30 min at 80 °C. MoS<sub>2</sub> was patterned into uniform rectangular channels of 2  $\mu\text{m}$  width by photolithography followed by a gentle O<sub>2</sub> plasma etch. The O<sub>2</sub> plasma etch was carried out at a power of 10 W in a chamber set to 20 mTorr pressure and 20 sccm O<sub>2</sub> flow rate for 60 s.

**Optical, Electrical, and AFM Characterizations.** Raman and PL data were taken in air, at room temperature, using a HORIBA Labram with a 532 nm excitation laser. All electrical measurements were at room temperature, in the dark and under vacuum (<10<sup>-5</sup> Torr) after an in situ vacuum anneal at 250 °C for 2 h, in a Janis ST-100 probe station. SEM was conducted using a Thermo Fisher Scientific Apreo S SEM equipped with a NiCol electron column operated in the immersion mode. AFM was conducted using a Bruker Dimension Icon AFM system operated in the tapping mode.

## AUTHOR INFORMATION

### Corresponding Author

**Eric Pop** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States; Department of Materials Science and Engineering and Precourt Institute for Energy, Stanford University, Stanford, California 94305, United States; [orcid.org/0000-0003-0436-8534](https://orcid.org/0000-0003-0436-8534); Email: [epop@stanford.edu](mailto:epop@stanford.edu)

### Authors

**Alvin Tang** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States; [orcid.org/0000-0003-4293-2917](https://orcid.org/0000-0003-4293-2917)

**Aravindh Kumar** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States; [orcid.org/0000-0002-4481-9910](https://orcid.org/0000-0002-4481-9910)

**Marc Jaikissoon** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States

**Krishna Saraswat** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States; Department of Materials Science and Engineering and Precourt Institute for Energy, Stanford University, Stanford, California 94305, United States

**H.-S. Philip Wong** – Department of Electrical Engineering, Stanford University, Stanford, California 94305, United States; Precourt Institute for Energy, Stanford University, Stanford, California 94305, United States

Complete contact information is available at:  
<https://pubs.acs.org/10.1021/acsami.1c06812>

### Author Contributions

A.T. and A.K. contributed equally. A.T., A.K., and E.P. conceived the experiments and wrote the manuscript with input from all authors. A.T. grew MoS<sub>2</sub> and performed the Raman and PL characterizations. A.K. and M.J. performed the SEM and AFM characterizations, respectively. A.K., M.J., and A.T. fabricated the devices. A.K. and A.T. performed the electrical measurements and data analysis.

### Notes

The authors declare no competing financial interest.

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